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Matched-field processing for leak localization in a 1 viscoelastic pipe: An experimental study 2 Xun Wang<sup>a</sup>, Jingrong Lin<sup>a</sup>, Alireza Keramat<sup>a</sup>, Mohamed S. Ghidaoui<sup>a</sup>, 3 Silvia Meniconi<sup>b</sup>, Bruno Brunone<sup>b</sup> 4 <sup>a</sup>Department of Civil and Environmental Engineering, Hong Kong University of Science 5 and Technology, Clear Water Bay, Hong Kong, China 6 <sup>b</sup>Department of Civil and Environmental Engineering, University of Perugia, 06125 7 Perugia, Italy 8

# 9 Abstract

This paper applies the matched-field processing (MFP) method to leakage 10 localization in a viscoelastic pipe. The viscoelasticity of pipe wall is included 11 in the governing equations of transient wave via the generalized Kelvin-Voigt 12 model and its effect is finally translated into a frequency-dependent wave 13 speed. Then, a leak is localized by MFP via a 1D search of leak location 14 along the pipe, independent of the leak size. Transient experiments with vis-15 coelastic pipe in the Water Engineering Laboratory at University of Perugia 16 and in the Water Resources Research Laboratory at Hong Kong University of 17 Science and Technology are studied. Experimental results demonstrate that 18 the inclusion of pipe wall viscoelasticity and using more frequencies (instead 19 of using only resonant frequencies) improve significantly the leak localization 20 accuracy. It is shown that the MFP leak localization is accurate even for a 21 small leak (the flow ratio of leak and main pipe is approximately 10%) in a 22 noisy environment: among 50 transient experiments, the maximum error of 23 MFP leak localization is only 1.14 m and in the other 49 experiments the 24 error is always lower than 1 m. 25

*Keywords:* transient wave, leakage localization, matched-field processing,
 pipe viscoelasticity, complex environment

# 28 1. Introduction

Leakage in water supply systems results in financial losses from wastage of water and health risks since leaks are potential entry points for contaminants

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during low pressure intrusion events [1]. Since 1980s, fluid transient-based 31 defect detection methodology has been used for leakage detection. It intro-32 duces hydraulic pressure waves, measures pressure response at specified loca-33 tion(s), and uses the information of reflection and damping due to leakages 34 to estimate their locations in water pipe systems. Specific methodological 35 examples of this approach are: (i) transient reflection-based method (TRM), 36 such as [2–7]; (ii) transient damping-based method (TDM) by [8]; (iii) fre-37 quency response-based method (FRM) by [9-25]; and (iv) inverse transient 38 analysis (ITA) method [26–29]. 30

Real water supply pipeline system is often a highly noisy environment due 40 to traffic, mechanical devices, turbulence, etc. While previous methods in the 41 literature do not theoretically or analytically study the effect of noise using 42 a probabilistic framework, a recent method, known as the matched-field pro-43 cessing (MFP) [21], estimate a leak based on the maximum signal-to-noise 44 ratio (SNR) meaning that the MFP method provides precise localization 45 estimates even in noisy environments. MFP is able to use all available fre-46 quencies, not just resonant frequencies, and does not need to identify which 47 frequencies are resonant frequencies, such that the leak estimation is more ro-48 bust. However, the MFP approach in [21] is based on a transient wave model 49 in elastic pipes; its availability in viscoelastic pipes has not been studied. 50

Viscoelastic pipes, such as polyvinyl chloride (PVC), polyethylene (PE), 51 and high-density polyethylene (HDPE), are ideally applicable in urban water 52 supply systems due to their excellence resilience, low cost, and convenience 53 in construction and maintenance. The viscoelastic effect of pipe deformation 54 during transient pressure behavior has been investigated [30-37]. It is shown 55 that the viscoelastic behavior changes the nature of transient wave. Leak 56 detection methods that can deal with pipe viscoelasticity have been proposed 57 [36, 38]. However, these methods do not consider the effect of noise in their 58 models and may thus not be applicable in a noisy environment. 59

In the present paper, the viscoelastic effect of pipe wall is included in the 60 transient model, it is found that it can be equivalently quantified by changing 61 the wave speed in the elastic case to be frequency-dependent. Then, MFP 62 can be applied for leakage detection in a viscoelastic pipe; its efficiency is 63 validated via experiments conducted in the Water Engineering Laboratory 64 at University of Perugia and in the Water Resources Research Laboratory 65 at Hong Kong University of Science and Technology. The accuracy of MFP 66 in pipeline leakage localization and additional gain of using more frequen-67 cies, instead of only the resonant frequencies, which have been numerically 68



Figure 1: Setup of the pipeline system.

<sup>69</sup> justified via numerical simulation and theoretical analysis in [21, 22], are ex-<sup>70</sup> perimentally illustrated in this paper. Experimental results show that the <sup>71</sup> inclusion of pipe viscoelasticity in the transient model and in the leak detec-<sup>72</sup> tion scheme significantly increases the leak localization accuracy.

The organization of this paper is as follows. It begins with a description of transient wave model in a viscoelastic pipe system in Section 2. Section 3 introduces the MFP method for leak localization. The setups, processing of data, and experimental results are all introduced in Section 4. Based on the experimental results of leak localization, Section 5 discusses the issues of physical model, measurement information, and leak localization algorithm. Finally, conclusions are drawn in Section 6.

# <sup>80</sup> 2. Transient wave in a viscoelastic pipe

#### <sup>81</sup> 2.1. Governing equations in the time domain

The pipeline configuration is illustrated in Figure 1. The upstream and 82 downstream ends of the single pipe locate at  $x = x^U = 0$  and  $x = x^D =$ 83 L, respectively. The pressure head  $h(x^U)$  at  $x^U$  is known and a transient 84 generator is placed at  $x^{D}$ . A leak is assumed whose location is denoted 85 by  $x^L$ , and  $z^L$ ,  $Q_0^L$  and  $H_0^L$  denote the elevation of the pipe at the leak, 86 the steady-state discharge and head at the leak, respectively. The lumped 87 leak parameter  $s^L = C^d A^L$  stands for the effective leak size, where  $C^d$  is 88 the discharge coefficient of the leak and  $A^{L}$  is the flow area of the leak 89 opening (orifice). The steady-state discharge of the leak is related to the 90 lumped leak parameter by  $Q_0^L = s^L \sqrt{2g(H_0^L - z^L)}$ , in which g is gravitational 91 acceleration. 92

For a viscoelastic pipe, the total strain  $\epsilon$  can be decomposed into an

# Nomenclature

q	discharge oscillation
h	pressure head oscillation
$x^L$	leak location
$z^L$	pipe elevation at leak
$s^L$	leak size
$Q_0^L, H_0^L$	steady-state discharge and head of leak
$Q_0$	steady-state discharge of main pipe
$x_m \ (m=1,\cdots,M)$	sensor coordinate
$\Delta \mathbf{h}$	head difference
n	measurement noise
$a_0$	wave speed in elastic pipes
$a_{ve}$	wave speed in viscoelastic pipes
A	area of pipeline
l	pipe length
D	pipe diameter
$\omega$	angular frequency
$\omega_{th}$	fundamental frequency
M	sensor number
K	frequency number
$\operatorname{FRF}$	frequency response function
FRM	frequency response-based method
HDPE	high-density polyethylene
K-V	Kelvin-Voigt
MFP	matched-field processing
MFP-E	MFP based on the model of elastic pipe $[21]$
MFP-VE	MFP based on the model of viscoelastic pipe
TDM	transient damping-based method
TRM	transient reflection-based method

instantaneous-elastic strain  $\epsilon_e$  and a retarded strain  $\epsilon_r$ , i.e.,

$$\epsilon = \epsilon_e + \epsilon_r. \tag{1}$$

Let  $\sigma(t)$  stand for dynamic stress at time t ( $\sigma(t) = 0$  for  $t \leq 0$ ),  $\epsilon_e$  and  $\epsilon_r$ read [39]:

97

$$\epsilon_e(t) = J_0 \sigma(t), \tag{2}$$

$$\epsilon_r(t) = \int_0^t \sigma(t-s) \frac{dJ}{ds} ds, \qquad (3)$$

<sup>98</sup> in which  $J_0 = J(0) = 1/E$  is the instantaneous creep-compliance representing <sup>99</sup> the immediate response of the material, E is the Young's modulus of elasticity <sup>100</sup> of pipe wall, and J(t) is the creep function.

The discharge and pressure head oscillations due to a fluid transient are denoted by q and h. The linearized unsteady-oscillatory continuity and momentum equations with time t and spatial coordinate  $x \in [x^U, x^L) \cup (x^L, x^D]$ [33, 34] are

$$\frac{\partial q}{\partial x} + \frac{gA}{a_0^2} \frac{\partial h}{\partial t} + 2(1-\nu^2)A \frac{\partial \epsilon_r}{\partial t} = 0$$
(4)

105 and

$$\frac{1}{gA}\frac{\partial q}{\partial t} + \frac{\partial h}{\partial x} + Rq = 0.$$
(5)

106 Here,

$$a_0 = \left(\rho\left(\frac{1}{\kappa} + (1-\nu^2)\frac{D}{Ee}\right)\right)^{-\frac{1}{2}} \tag{6}$$

is the elastic component of wave speed,  $\kappa$  and  $\rho$  are the bulk modulus and 107 density of water,  $\nu$  is the Poisson's ratio, D is the pipe diameter, and e 108 is the pipe wall thickness. Furthermore, A is the area of pipeline, R is the 109 steady-state resistance term being  $R = (f_{DW}Q_0)/(gDA^2)$  for turbulent flows, 110  $f_{DW}$  is the Darcy-Weisbach friction factor,  $Q_0$  is the steady-state discharge 111 in the pipe. The last term of left hand side of Eq. (4) is due to pipe wall 112 viscoelasticity; it equals to 0 in the elastic case since the creep function J(t) is 113 time-independent. Furthermore,  $\epsilon_r$  here is the retarded circumferential strain 114 and the influence of axial pipe velocity on Eqs. (4) and (5) is neglected. 115 Inserting Eq. (3) into Eq. (4) gives 116

$$\frac{\partial q}{\partial x} + \frac{gA}{a_0^2}\frac{\partial h}{\partial t} + 2(1-\nu^2)A\frac{\partial}{\partial t}\int_0^t \sigma(t-s)\frac{dJ}{ds}ds = 0.$$
 (7)

<sup>117</sup> Considering the force balance for the stress in the pipe wall and the pressure<sup>118</sup> in fluid, i.e.,

$$2\sigma e = pD, \tag{8}$$

<sup>119</sup> where  $p = \rho g h$  is the dynamic pressure, we have

$$\sigma(t) = \frac{\rho g D}{2e} h(t). \tag{9}$$

 $_{120}$  Therefore, Eq. (7) becomes

$$\frac{\partial q}{\partial x} + \frac{gA}{a_0^2}\frac{\partial h}{\partial t} + (1-\nu^2)\frac{\rho gDA}{e}\frac{\partial}{\partial t}\int_0^t h(t-s)\frac{dJ}{ds}ds = 0.$$
(10)

- 2.2. Governing equations in the frequency domain and wave speed calibration
   for pipe wall viscoelasticity
- Taking a Fourier transform of Eq. (10), it becomes

$$\frac{\partial q}{\partial x} + \frac{gA}{a_0^2} i\omega h + (1 - \nu^2) \frac{\rho g DA}{e} \mathcal{F}\left(\frac{\partial}{\partial t} \int_0^t h(t - s) \frac{dJ}{ds} ds\right) = 0, \qquad (11)$$

where  $\mathcal{F}(\cdot)$  stands for the operation of Fourier transform. The last term of Eq. (11) can be simplified by the Leibniz's rule for differentiation under the integral sign, the properties of Fourier transform and convolution, and h(t) = 0 for  $t \leq 0$ , as:

$$\mathcal{F}\left(\frac{\partial}{\partial t}\int_{0}^{t}h(t-s)\frac{dJ}{ds}ds\right)$$

$$= \mathcal{F}\left(h(0)\frac{dJ}{dt} + \int_{0}^{t}\frac{\partial h}{\partial t}(t-s)\frac{dJ}{ds}ds\right) = \mathcal{F}\left(\int_{0}^{t}\frac{\partial h}{\partial t}(t-s)\frac{dJ}{ds}ds\right)$$

$$= \mathcal{F}\left(\frac{\partial h}{\partial t}\right)\mathcal{F}\left(\frac{dJ}{dt}\right) = i\omega h\mathcal{F}\left(\frac{dJ}{dt}\right).$$
(12)

Here, the creep function J(t) is assumed to follow the generalized Kelvin-Voigt (K-V) model [31]:

$$J(t) = J_0 + \sum_{j=1}^{N_{kv}} J_j (1 - \exp(-t/\tau_j)), \qquad (13)$$

where  $N_{kv}$  is the truncated order,  $J_j$  and  $\tau_j$  are coefficients of the K-V model. Therefore, we have

$$\frac{dJ}{dt} = \sum_{j=1}^{N_{kv}} \frac{J_j}{\tau_j} \exp(-t/\tau_j) \tag{14}$$

132 and

$$\mathcal{F}\left(\frac{dJ}{dt}\right) = \sum_{j=1}^{N_{kv}} \frac{J_j}{1 + \mathrm{i}\omega\tau_j}.$$
(15)

<sup>133</sup> By Eqs. (12) and (15), Eq. (11) becomes

$$\frac{\partial q}{\partial x} + \frac{gA}{a_{ve}^2} i\omega h = 0, \qquad (16)$$

134 where

$$a_{ve} = \left(\rho\left(\frac{1}{\kappa} + (1-\nu^2)\frac{D}{e}\left(J_0 + \sum_{j=1}^{N_{kv}} \frac{J_j}{1+i\omega\tau_j}\right)\right)\right)^{-\frac{1}{2}}.$$
 (17)

<sup>135</sup> Furthermore, taking Fourier transform of Eq. (5), we obtain

$$\frac{\partial h}{\partial x} + \left(\frac{\mathrm{i}\omega}{gA} + R\right)q = 0. \tag{18}$$

Eqs. (16) and (18) are respectively the continuity and momentum equations in the frequency domain. In the considered viscoelastic case where the pipe motion is assumed to be purely radial, the momentum equation Eq. (18) is same as the elastic case; the continuity equation changes but it is equivalent to replace the constant wave speed in the elastic case ( $a_0$  in Eq. (6)) by the frequency-dependent  $a_{ve}$  in Eq. (17).

#### 142 2.3. Boundary conditions and data model

Solving Eqs. (16) and (18) with boundary conditions of the discharge  $q(x^U)$  and head  $h(x^U)$  at  $x^U$ , for the case of  $x_m < x^L$ , the discharge and head (frequency response [40]) at a measurement station  $x_m$  can be computed as [41, 42]:

$$\begin{pmatrix} q(x_m) \\ h(x_m) \end{pmatrix} = M^{NL}(x_m) \begin{pmatrix} q(x^U) \\ h(x^U) \end{pmatrix},$$
(19)

147 in which

$$M^{NL}(x) = \begin{pmatrix} \cosh(\mu x) & -\frac{1}{Z}\sinh(\mu x) \\ -Z\sinh(\mu x) & \cosh(\mu x) \end{pmatrix}$$
(20)

148 is the field matrix,

$$Z(\omega) = \mu a_{ve}^2(\omega) / (i\omega gA)$$
(21)

<sup>149</sup> is the characteristic impedance,

$$\mu(\omega) = a_{ve}^{-1}(\omega)\sqrt{-\omega^2 + igA\omega R}$$
(22)

is the propagation function. In the case of  $x_m > x^L$ , by the head and mass conservation condition across the leak:

$$h(x^{L-}) = h(x^{L+}) = h(x^{L});$$
 (23)

152

$$q(x^{L-}) = q(x^{L+}) + q(x^{L}) = q(x^{L+}) + \frac{Q_0^L}{2(H_0^L - z^L)}h(x^L), \qquad (24)$$

<sup>153</sup> in which  $x^{L-}$  and  $x^{L+}$  represent respectively just upstream and downstream <sup>154</sup> of  $x^{L}$ , the discharge and head at a measurement station  $x_m$  has the form <sup>155</sup> [41, 42]:

$$\begin{pmatrix} q(x_m) \\ h(x_m) \end{pmatrix} = M^{NL}(x_m - x^L) \begin{pmatrix} 1 & -\frac{Q_0^L}{2(H_0^L - z^L)} \\ 0 & 1 \end{pmatrix} M^{NL}(x^L) \begin{pmatrix} q(x^U) \\ h(x^U) \end{pmatrix}.$$
(25)

The transfer matrix on the right hand side of Eq. (25) can be further simplified as a form of variable separation of  $x^L$  and  $s^L$  [21, 22]; the head at  $x = x_m$  for a given angular frequency  $\omega_k$  is [21]:

$$h(\omega_k, x_m) = h^{NL}(\omega_k, x_m) + s^L G(\omega_k, x^L, x_m),$$
(26)

159 wherein

$$h^{NL}(\omega_k, x_m) = -Z_k \sinh(\mu_k x_m) q(x^U, \omega_k) + \cosh(\mu_k x_m) h(x^U, \omega_k)$$
(27)

160 and

$$= \begin{cases} G(\omega_k, x^L, x_m) \\ -\frac{\sqrt{g}Z_k \sinh(\mu_k(x_m - x^L))}{\sqrt{2(H_0^L - z^L)}} \left( Z_k \sinh(\mu_k x^L) q(x^U, \omega_k) - \cosh(\mu_k x^L) h(x^U, \omega_k) \right), & x_m > x^L \\ 0, & x_m \le x^L \end{cases}$$

$$(28)$$

Assume that the measured heads at the frequency  $\omega_k$   $(k = 1 \cdots, K)$  and at the location  $x_m$   $(m = 1 \cdots, M)$  are available and they are assumed to be contaminated by a noise  $n_{mk}$ , i.e.,

$$h(\omega_k, x_m) = h^{NL}(\omega_k, x_m) + s^L G(\omega_k, x^L, x_m) + n_{mk}.$$
 (29)

164 By denoting

$$\Delta \mathbf{h} = (\Delta h_{11}, \cdots, \Delta h_{1K}, \cdots, \Delta h_{M1}, \cdots, \Delta h_{MK})^{\mathsf{T}}, \qquad (30)$$

165 where

$$\Delta h_{mk} = h(\omega_k, x_m) - h^{NL}(\omega_k, x_m), \qquad (31)$$

<sup>166</sup> 
$$\mathbf{G}(x^{L}) = (G(\omega_{1}, x^{L}, x_{1}), \cdots, G(\omega_{K}, x^{L}, x_{1}), \cdots, G(\omega_{1}, x^{L}, x_{M}), \cdots, G(\omega_{K}, x^{L}, x_{M}))^{\top},$$
(32)

167 and

$$\mathbf{n} = (n_{11}, \cdots, n_{1K}, \cdots, n_{M1}, \cdots, n_{MK})^{\mathsf{T}},$$
(33)

168 we finally have

$$\Delta \mathbf{h} = s^L \mathbf{G}(x^L) + \mathbf{n}. \tag{34}$$

In this paper,  $\Delta \mathbf{h}$  is the data for leakage localization algorithm.

Note that in Eq. (28), the boundary conditions at the upstream node 170  $q(x^U)$  and  $h(x^U)$  are assumed to be known. Here,  $h(x^U)$  at the upstream 171 boundary  $x^U$  is known, for example  $h(x^U)$  can be reasonably assumed to be 172  $h(x^U) = 0$  if the upstream is connected to a reservoir. The discharge  $q(x^U)$ 173 can be estimated if a transducer near the upstream boundary, whose location 174 is denoted by  $x_0$ , is available [43, 44]. Assuming there is no leak between  $x^U$ 175 and  $x_0$  and using the pressure head measurement  $h(x_0)$  at  $x_0$ , the discharge 176  $q(x^U)$  at the frequency  $\omega_k$  can be estimated [22] by 177

$$\hat{q}(x^{U},\omega_{k}) = \frac{\cosh(\mu_{k}(x_{0}-x^{U}))h(x^{U},\omega_{k}) - h(x_{0},\omega_{k})}{Z_{k}\sinh(\mu_{k}(x_{0}-x^{U}))}.$$
(35)

# 178 3. Leak localization using the matched-field processing method

In this section, the leakage localization problem is solved using the MFP method. The noise vector **n** is assumed to follow a zero-mean Gaussian distribution  $\mathcal{N}(\mathbf{0}, \sigma^2 \mathbf{I}_{MK})$ , where  $\mathbf{I}_{MK}$  is a MK-dimensional identity matrix. The leakage location can be estimated using MFP via [21]:

$$\widehat{x^{L}} = \arg \max_{x^{L} \in (0,l)} \frac{\Delta \mathbf{h}^{H} \mathbf{G}(x^{L}) \mathbf{G}^{H}(x^{L}) \Delta \mathbf{h}}{\mathbf{G}^{H}(x^{L}) \mathbf{G}(x^{L})}.$$
(36)

Remark that the random noise is assumed to be Gaussian distributed. Since 183 frequency-domain pressure is obtained from discrete Fourier transform, which 184 is a linear combination of measured time-domain pressure signals, thus the 185 frequency-domain noise can be approximately as Gaussian distributed ac-186 cording to the Law of Large Number. A recent experimental investigation 187 [45] also shows that, in a pipe with flow, the Gaussian assumption of noise 188 distribution is reasonable. Furthermore, uncorrelated white noise (the co-189 variance matrix is  $\sigma^2 \mathbf{I}_{MK}$ ) is assumed here; if the noise is non-white, a data 190 transformation technique for noise whitening [21] can be applied before im-191 plementing MFP such that the white noise assumption still holds and the 192 leakage localization in Eq. (36) can still be used without changing its struc-193 ture. This noise-whitening technique can improve the leak localization result 194 when noise level is high according to the numerical results in [21]. 195

Note that Eq. (36) implies that a leak can be localized by a 1D search of leak location along the pipe, independent of its leak size. Physically, this is possible because the leak location determines the shape of FRF while the leak size only proportionally changes the magnitude of FRF at different frequencies [15, 46]. Essentially, MFP uses only the shape of FRF [21] such that the influence of leak size can be excluded.

Finally, the MFP algorithm of leakage localization in a viscoelastic pipe is summarized in Algorithm 1.

#### **4.** Experimental results

# 205 4.1. Experiments at University of Perugia

In this section, the MFP method for leak localization is tested using the water hammer experimental data obtained from the Water Engineering Laboratory of University of Perugia.

#### 209 4.1.1. Experimental setup

The experimental setup is shown in Figure 2. A HDPE pipe with length 210 l = 166.28 m and diameter D = 0.0933 m is used. Transient wave is generated 211 by a rapid and full closure of the downstream value; the time duration of value 212 closure is approximately 0.1 s. Pressures are measured by three transducers 213 located at  $x_0 = 27.7$  m,  $x_1 = 68.27$  m, and  $x_2 = 166.28$  m, respectively. The 214 average steady-state pressure head (pre-transient) at the upstream reservoir 215 and at the valve are respectively 18.28 m and 17.09 m. The leak location 216 is  $x^L = 60.84$  m and the effective leak size is  $C^d A^L = 6.8 \times 10^{-5}$  m<sup>2</sup>. By 217

Algorithm 1 Localization of leakage in a viscoelastic pipe using MFP

- 1. Select K frequencies  $\omega_1, \dots, \omega_K$ . The maximum frequency  $\omega_K$  should not be higher than the maximum probing frequency. It is also suggested to use all available frequencies that are lower than the maximum probing frequency. The maximum probing (angular) frequency can be approximately computed from  $2\pi/T_c$  ( $T_c$  is the time duration of valve closure) or observed from the spectrum of measured signal.
- 2. Compute  $a_{ve}$  from Eq. (17) for each selected frequency  $\omega_k$   $(k = 1, \dots, K)$ , in which the viscoelastic coefficients  $N_{kv}$ ,  $J_j$  and  $\tau_j$   $(j = 1, \dots, N_{kv})$  of the generalized K-V model are used and obtained prior to the transient test with leak.
- 3. Estimate the boundary condition  $\hat{q}(x^U, \omega_k)$  for each selected frequency  $\omega_k$  from Eq. (35) using the pressure head measurements from  $x_0$ .
- 4. Calculate  $h^{NL}(\omega_k, x_m)$  via Eq. (27) and use the head differences  $\Delta \mathbf{h}$  as the data, which includes pressure heads from K frequencies and M sensors  $(x_1, \dots, x_M)$ .
- 5. Plot the objective function in Eq. (36):

$$|B|^{2} = \frac{\Delta \mathbf{h}^{H} \mathbf{G}(x^{L}) \mathbf{G}^{H}(x^{L}) \Delta \mathbf{h}}{\mathbf{G}^{H}(x^{L}) \mathbf{G}(x^{L})}.$$
(37)

with respect to  $x^L$  and retain  $x^L$  corresponding to maximum  $|B|^2$  as leak location estimate.





Figure 2: (a) Photo and (b) sketch map of pipe transient experiment in the Water Engineering Laboratory at University of Perugia.

computing the travel time of the wave between the transducers, the wave speed is estimated as  $a_0 = 375$  m/s. The Darcy-Weisbach friction factor is also computed from the measured head loss as  $f_{DW} = 0.0233$ .

## 221 4.1.2. Post-processing of measurements

The measured heads in time from the three sensors are shown in Figure 3. Their system frequency response functions (FRFs) are obtained as follows.

• Computing the FRF requires both the output from the system (the 224 measured pressure head of transient wave in Figure 3) and the input 225 signal sent to the system. The latter is computed by selecting only 226 the first step rise (from the last steady-state point to the maximum 227 of the first jump) of the measurement and keeping it constant after 228 the maximum point [47], since the input signal is generated by the full 229 closure of valve. Figure 4 (a) takes the measurement from  $x_1$  as an 230 example which displays the measured head and the computed input 231 signal. 232

- Note that these step time-series signals have infinite energy in theory 233 (their integral from  $-\infty$  to  $+\infty$  equals to infinity), such that their 234 spectra cannot be produced by a conventional Fourier transform [48]. 235 Therefore, in order to change the signals into a finite energy form, 236 each step signal is modified to a pulse-type signal by computing the 237 difference between the original signal and its delay. The delay of the 238 measurement or the input signal is computed by delaying the original 239 signal with a time lag being longer than the time from the last steady-240 state point to the maximum of the first jump and being shorter than the 241 time for the first reflection to arrive at the measure station [47]. The 242 accuracy of this correction procedure does not depend on the time lag 243 factor if it is in this range. Figure 4 (b) and (c) show the original signals 244 and their delays of the measurement and the input, while Figure 4 (d) 245 displays the differences (impulse-type signals) of the measured head 246 and the input signal. According to the additive and distributive nature 247 of time invariant linear systems, the frequency response of the system 248 remains unchanged as a result of this operation (cf. Eq. (5.21) in [47]). 249
- 250 Т 251 О

• Then, conventional Fourier transform is applied to the impulse-type output and input signals, which gives the corresponding frequency do-



Figure 3: Measured pressure in the time domain from the three transducers.

252 253

main signals.	Finally, the	e system FF	RF is obtained	by their quotie	nt.
The FRFs obt	tained from	the three se	ensors are show	n in Figure 5.	

$\nu = 0.43$	$\kappa = 2.1 \times 10^9$ Pa	$ ho = 10^3 \ \mathrm{kg/m^3}$
$D = 93.3 \times 10^{-3} \text{ m}$	$e = 7.5 \times 10^{-3} \text{ m}$	$J_0 = 0.68 \times 10^{-9} \text{ Pa}^{-1}$
$J_1 = 0.951 \times 10^{-10} \text{ Pa}^{-1}$	$\tau_1 = 0.05 \ s$	$J_2 = 1.065 \times 10^{-10} \text{ Pa}^{-1}$
	$J_3 = 0.815 \times 10^{-10} \text{ Pa}^{-1}$	$ au_3 = 1.5 \ s$

254 4.1.3. Leak localization results

Table 1: Coefficients for pipe wall viscoelasticity in the experiments at University of Perugia.

As already stated, the wave speed is estimated by travel time between two transducers and this leads to  $a_0 = 375$  m/s. The calibrated frequencydependent wave speed  $a_{ve}$  is obtained from Eq. (17), where the truncated order  $N_{kv} = 3$  and the corresponding coefficients  $J_j$  and  $\tau_j$ , j = 1, 2, 3, are obtained (shown in Table 1) via a calibration of viscoelastic parameters of



Figure 4: Post-processing of measurement from the transducer at  $x_1$ : (a) the measurement and the computed input signal; (b) the measurement and its time delay; (c) the input signal and its time delay; (d) difference of the measurement and its delay and the difference of the input signal and its delay.



Figure 5: Frequency response function of measurements from the three transducer.

pipe material [49]. More specifically, a transient test without leak is done prior to the leaking test and the viscoelastic parameters are calibrated such that the transient wave model (without leak but  $J_j$  and  $\tau_j$  are free parameters) is closest to the measured data.

In the following, the leak is localized using MFP. The leak detection results with and without the viscoelastic terms are both shown. To avoid confusion, we denote MFP-E as the MFP method based on the model of elastic pipe that does not include the viscoelastic terms (i.e., the method in [21]) and MFP-VE as the viscoelastic version of MFP proposed in the present paper. In MFP-E, the wave speed is  $a_0 = 375$  m/s.

The upstream discharge  $q(x^{U})$ , which appears in both  $h^{NL}$  and G in 270 Eq. (34), is estimated via Eq. (35), in which  $h(x_0)$  is the frequency response 271 from the upstream transducer  $x_0 = 27.7$  m. The choice of frequencies is an 272 important issue for MFP. First, the maximum frequency included in the data 273  $\Delta \mathbf{h}$  (or bandwidth) should be decided. In Figure 5, the shape of FRF can 274 be clearly seen for low frequencies, but the signal of transient wave becomes 275 weak as the frequency increases. As a matter of fact, it can be observed 276 from Figure 5 that the data of FRF are not reliable when  $\omega/\omega_{th} > 15$  in 277 this experiment and is mainly comprised of noise. Second, the number of 278 used frequency (in the given frequency range) also affects the localization 279 results. It has been shown in [21] that with a given bandwidth using more 280 frequencies increases the robustness of leak localization, particularly in a 281 noisy environment. 282

Figure 6 shows the leak localization results using MFP. Here, the used 283 angular frequencies are denoted by  $\{\omega : \omega_{th}, \omega_{th} + \Delta \omega, \omega_{th} + 2\Delta \omega, \cdots, n\omega_{th}\}$ . 284 Figure 6 (a) and (b) displays the results for n = 13 (i.e., seven peaks in 285 FRFs) and  $\Delta \omega = \omega_{th}/7$  (i.e., in total 85 frequencies are used) using MFP-E 286 and MFP-VE, respectively. It is clear that the model without viscoelasticity 287 leads to a bias of leak localization (the error is 1.74 m) and the model 288 with viscoelasticity improves the precision (the error is 0.66 m). This illus-289 trates the importance of including viscoelasticity in the detection scheme. 290 Figure 6 (c) and (d) shows the corresponding results with resonant frequen-291 cies only. In the case without including pipe viscoelasticity, a higher peak of 292 MFP objective function appears near the downstream end of the pipe which 293 results in a wrong estimate of leak location. The result is better when the 294 viscoelasticity is considered but the error (1.76 m) is still larger than using 295 more frequencies (0.66 m in Figure 6 (b)). 296

The localization error with different frequency bandwidth  $(n = 5, 7, \dots, 17)$ 

and frequency step size  $(\Delta \omega = \omega_{th}/7, \omega_{th}/2, \omega_{th}, 2\omega_{th})$  are shown in Figure 7. 298 These results clearly show that MFP-VE improves the leak localization accu-299 racy. Furthermore, this figure justifies that using more frequencies in a given 300 range decreases the localization error. In the cases where  $\Delta \omega = \omega_{th}/7, \omega_{th}/2, \omega_{th}/2,$ 301 the error is always less than 5 m (labeled by the dotted lines in Figure 7) for 302 all the frequency ranges except n = 5 where too little information is available. 303 By contrast, only using resonant (or even with anti-resonant frequencies to-304 gether) obviously increases the error; including viscoelasticity becomes more 305 important and the choice of frequency range must be careful. The working 306 frequency (here, "working" means the localization error is less than 5 m) for 307 the four cases of frequency densities is listed in the second column of Table 2. 308 This table also displays the average error for each  $\Delta \omega$  in the corresponding 309 working range of frequency, which illustrates again the importance of includ-310 ing viscoelasticity in the leak detection scheme and using more frequencies. 311

$\Delta \omega$	n	average error [m]				
		MFP-E	MFP-VE			
$2\omega_{th}$	7 - 11	1.39	0.89			
$\omega_{th}$	7 - 13	1.21	0.99			
$\omega_{th}/2$	7 - 17	1.04	0.72			
$\omega_{th}/7$	7 - 17	0.96	0.58			

Table 2: The chosen maximum frequency such that the leak localization error less than 5 m (represented by n in the second column) and the average error of leak localization in the corresponding frequency range with various frequency spacing  $\Delta \omega = 2\omega_{th}, \omega_{th}, \omega_{th}/2, \omega_{th}/7$ .



Figure 6: Leak localization using MFP with different frequency spacing: (a,b)  $\Delta \omega = \omega_{th}/7$  (i.e., in total 85 frequencies are used) and (c,d)  $\Delta \omega = 2\omega_{th}$  (resonant frequencies only; in total 7 frequencies). In subfigures (a,c) and subfigures (b,d), the results are obtained with MFP-E and MFP-VE, respectively. The crosses and dash lines represent the sensor and actual leak locations, respectively.



Figure 7: Leak localization error with different frequency spacing (a)  $\Delta \omega = 2\omega_{th}$  (resonant frequencies only); (b)  $\Delta \omega = \omega_{th}$  (resonant and anti-resonant frequencies); (c)  $\Delta \omega = \omega_{th}/2$ ; and (d)  $\Delta \omega = \omega_{th}/7$ . The solid and dash lines respectively stand for MFP-E and MFP-VE.

# 312 4.2. Experiments at Hong Kong University of Science and Technology

This section applies the MFP leak localization method to the experimental results obtained from the newly-built water pipe system at the Water Resources Research Laboratory at Hong Kong University of Science and Technology.

## 317 4.2.1. Experimental setup and measurements

The setup of HDPE pipeline system is shown in Figure 8. The pipe 318 length is l = 144 m and its diameter is D = 0.0792 m. A pump (in Figure 8) 319 (a) middle), instead of the surge tank used in the experiments conducted at 320 University of Perugia, is connected to the upstream of pipe. A downstream 321 valve is used to generate transient waves; the time closure is approximately 322 0.05 s. Three hydrophones at  $x_0 = 36.92$  m,  $x_1 = 141.43$  m and  $x_2 = 141.43$  m and  $x_2 = 141.43$  m and  $x_3 = 141.43$  m and  $x_4 = 141.43$  m and  $x_5 = 141.43$  m and  $x_8 = 141.43$  m and  $x_8$ 323 122.25 m are set in the pipe to measure pressure. A leak locates at  $x^{L}$  = 324 45.58 m (in Figure 8 (a) right); the leak size can be changed via different 325 valve settings. Three different leak sizes are tested, as shown in Figure 9, 326 because leak size is crucial in terms of leak detectability [50]. The three 327 tests correspond to the following steady-state flow ratio of leak and main 328 pipe:  $Q_0^L/Q_0 = 40\%, 20\%, 10\%$ . The leak flow  $Q_0^L$  and the pipe flow  $Q_0$  are 329 measured from two flow meters just upstream and just downstream of the 330 leak. 331

Figure 10 (a) shows three different measured pressure head signals (i.e., 332 from three different water hammer experiments) at  $x_1 = 141.43$  m, where 333  $Q_0^L/Q_0 = 40\%$ . It can be seen from this figure that the measured signal 334 is much more noisy and aleatory, and has more uncertainties than those in 335 Figure 3 (experiments at University of Perugia). A similar phenomenon can 336 be observed by the corresponding FRF in Figure 10 (b), which is also much 337 more noisy than those in Figure 5. This may be partially due to the pump, 338 which itself generates much noise. The presence of the noise and uncertainty 339 makes the leak localization more challenging: 340

- In Measurement 2 in Figure 10 (a), the pressure drop due to leak is not obvious even if the leak is large  $(Q_0^L/Q_0 = 40\%)$ . This unclear reflection due to leak may mean that some reflection-based leak detection methods fail.
- In Measurement 3 in Figure 10 (a), a very strong noise can be seen in the first period of time signal, which may disturb some signal processing techniques.



Figure 8: (a) Photos (pipe (left), pump (middle), and leak (right)) and (b) sketch map of pipe transient experiment in the Water Resources Research Laboratory at Hong Kong University of Science and Technology.



Figure 9: Photos of three different leak sizes in the HKUST experiments. The flow rate ratio between the main and the leak is approximately 40% (left), 20% (middle), and 10% (right)

• The FRFs in Figure 10 (b) are noisy and the resonant frequencies are not so clear as in the experiments at University of Perugia (cf. Figure 5). This poses a challenge for methods based on resonant frequencies.

Moreover, the MFP method needs  $\hat{q}(x^U)$  which is estimated via Eq. (35). Here, it is assumed  $h(x^U) = 0$  in the computation of Eq. (35). However, unlike the experiments in Section 4.1 where a reservoir is connected to the pipe upstream to keep the pressure head constant, in the experiments in this section a pump is used such that the assumption  $h(x^U) = 0$  could result in uncertainties in the model.

#### 358 4.2.2. Leak localization results

Figure 11 shows the MFP leak localization results using the signals in 359 Figure 10. Here, the frequencies used are all available frequencies from  $\omega_{th}$ 360 to  $17\omega_{th}$ . The MFP without viscoelastic calibration (MFP-E) results in bad 361 results; in Figure 11 (a,c,e), many peaks with almost equal height appear in 362 the MFP objective function. However, the leak can be accurately localized 363 by the proposed MFP with consideration of viscoelasticity (MFP-VE), where 364 the coefficients of viscoelasticity are displayed in Table 3, for all the three 365 measurements (the error is 0.14 m, 0.96 m and 0.26 m, respectively), which 366 illustrate the robustness of the proposed method with respect to noise and 367 uncertainties. 368

$\nu = 0.43$	$\kappa = 2.1 \times 10^9$ Pa	$\rho = 10^3 \text{ kg/m}^3$
$D = 79.2 \times 10^{-3} \text{ m}$	$e = 5.4 \times 10^{-3} \mathrm{~m}$	$J_0 = 4.8 \times 10^{-10} \text{ Pa}^{-1}$
$J_1 = 1.957 \times 10^{-10} \text{ Pa}^{-1}$	$\tau_1 = 0.038 \text{ s}$	$J_2 = 1.0954 \times 10^{-10} \text{ Pa}^{-1}$
$ au_2 = 0.6 \ s$	$J_3 = 9.05 \times 10^{-12} \text{ Pa}^{-1}$	$ au_3 = 1.5 \ s$

Table 3: Coefficients for pipe wall viscoelasticity in the experiments at Hong Kong University of Science and Technology.

<sup>369</sup> Due to the high level of noise and uncertainty, for each leak size  $(Q_0^L/Q_0 = 40\%, 20\%, 10\%)$  the experiment is repeated 50 times and the leak localization <sup>371</sup> results are statistically analyzed. Here, we compare MFP-VE and MFP-E <sup>372</sup> with only resonant frequencies and with more frequencies, as well as three <sup>373</sup> representative methods in the literature:



Figure 10: Three different measurements obtained from the sensor  $x_1 = 141.43$  m in the time domain (a) and in the frequency domain (b).



Figure 11: Leak localization using MFP-E (a,c,d) and MFP-VE (b,d,f). The three measurements shown in Figure 10 are respectively used.

- (a): MFP-VE (the method proposed in this paper) with all available frequencies  $\omega \in [\omega_{th}, 17\omega_{th}]$ .
- (b): MFP-VE (the method proposed in this paper) with only resonant frequencies, i.e.,  $\omega \in \{\omega_{th}, 3\omega_{th}, \cdots, 17\omega_{th}\}.$
- (c): MFP-E (the method in [21]) with all available frequencies  $\omega \in [\omega_{th}, 17\omega_{th}]$ .
- (d): MFP-E (the method in [21]) with only resonant frequencies, i.e.,  $\omega \in \{\omega_{th}, 3\omega_{th}, \cdots, 17\omega_{th}\}.$
- (e): The wavelet analysis method [5] (a representative of TRM) where the mother wavelet is the type Daubechies of order 1 (db1). It analyzes the drop due to leak in the first period of time signal from the hydrophone at  $x_1$  and retains the largest drop as leak estimate.
- (f): The FRF peak pattern method [14, 16] (a representative of FRM; more specifically, a representative of resonant frequency based-method) which uses the peak frequencies  $\omega \in \{\omega_{th}, 3\omega_{th}, \cdots, 17\omega_{th}\}$  of FRF of the hydrophone at  $x_1$  as data for leak localization.
- (g): The inverse transient analysis (ITA) method [26] with consider-390 ation of viscoelastic effect [38]. Here, 51 potential leaks uniformally 391 distributed in the pipe, i.e.,  $\{x : x = nl/50, n = 0, \dots, 50\}$ , and the 392 sizes of the 51 potential leaks are estimated by solving a 51-parameter 393 optimization problem. The one with highest leak size estimate is re-394 tained as leak location estimate. The measurements from all the three 395 sensors are used. The information of actual leak location is used for the 396 initialization of optimization (although it is not available in practice): 397 for all the three cases of  $Q_0^L/Q_0$ , the initial leak size at the potential 398 leak closest to the actual leak is  $3.4 \times 10^{-5}$  m<sup>-2</sup>, while at the other 399 potential leaks the initial sizes are all 0. 400

Figures 12, 13, and 14 show the box plots of absolute error of leak localization, denoted by  $|e| = |\widehat{x^L} - x^L|$ , in the cases of three leak sizes. Tables 4, 5, and 6 display the corresponding statistics |e|, including mean, standard deviation, maximum, minimum, and percentiles of |e|. The results illustrate that MFP-VE with all frequencies is very accurate in all the three cases of

leak sizes. For example, for the smallest leak where  $Q_0^L/Q_0 = 10\%$  the aver-406 age error is only 0.56 m, the maximum error is 1.16 m, and in 49 over 50 tests 407 the leak localization error is less than 1 m. The average error is lower for 408 the larger leaks. Both considering viscoelasticity and using more frequencies 409 are essential for accuracy of MFP leak localization: the methods (b), (c) and 410 (d) all have much higher error than (a). The methods (e), (f) and (g) also 411 have relatively large error of leak localization. It is remarkable that for the 412 FRF peak pattern method (f), due to only nine FRF peaks are available, the 413 resolution of leak localization is actually very low that leak location estimate 414 can only be chosen from nine candidates along the pipe. Similarly, the res-415 olution of ITA method (g) is decided by the distribution of potential leaks; 416 here, the resolution of (g) is l/100 = 1.44 m because the leak is selected from 417  $\{x: x = nl/50, n = 0, \dots, 50\}$  (51 potential leaks are assumed in this case). 418 In practice, excavation cost is proportional to range of possible leak lo-419 cation, therefore leak localization with accuracy less than a threshold is of 420 importance. Successful rate of leak localization in the sense of absolute error 421 |e| less than 1 m, 3 m, 5 m and 10 m, computed from the 50 water hammer 422

experiments, is shown in Table 7. Again, MFP-VE with all frequencies (a) outperforms other methods. The wavelet method (e), which does not need many pipe parameters such as viscoelasticity coefficients, has error less than 5 m in most cases which is also acceptable in this sense.



Figure 12: Box plot of absolute error of leak localization obtained from 50 water hammer experiments. The methods used for leak localization is (a) MFP-VE with all frequencies; (b) MFP-VE with only resonant frequencies; (c) MFP-E with all frequencies; (d) MFP-E with only resonant frequencies; (e) wavelet transform method; (f) FRF peak pattern method; (g) ITA method. The ratio  $Q_0^L/Q_0$  of leak flow and main flow is 40%.



Figure 13: Box plot of absolute error of leak localization obtained from 50 water hammer experiments. The methods used for leak localization is (a) MFP-VE with all frequencies; (b) MFP-VE with only resonant frequencies; (c) MFP-E with all frequencies; (d) MFP-E with only resonant frequencies; (e) wavelet transform method; (f) FRF peak pattern method; (g) ITA method. The ratio  $Q_0^L/Q_0$  of leak flow and main flow is 20%.



Figure 14: Box plot of absolute error of leak localization obtained from 50 water hammer experiments. The methods used for leak localization is (a) MFP-VE with all frequencies; (b) MFP-VE with only resonant frequencies; (c) MFP-E with all frequencies; (d) MFP-E with only resonant frequencies; (e) wavelet transform method; (f) FRF peak pattern method; (g) ITA method. The ratio  $Q_0^L/Q_0$  of leak flow and main flow is 10%.

<i>e</i>	(a)	(b)	(c)	(d)	(e)	(f)	(g)
Mean [m]	0.29	19.49	20.17	73.16	8.89	13.43	6.47
Std. [m]	0.27	16.72	26.21	24.53	17.51	8.19	3.17
Max [m]	1.24	35.54	70.14	91.54	74.95	26.42	15.26
Min [m]	0.04	0.14	4.36	4.06	0.07	2.38	0.31
5% Percentile [m]	0.04	0.14	4.76	4.36	0.21	2.38	0.5
25% Percentile [m]	0.14	0.76	5.06	72.64	0.82	2.38	4.49
50% Percentile [m]	0.16	33.19	5.26	73.09	1.74	16.78	6.95
75% Percentile [m]	0.44	34.44	24.84 m	90.64	4.77	73.16	8.14
95% Percentile [m]	0.96	35.04	69.64 m	91.24	65.18	26.42	12.02

Table 4: Statistics of absolute error of leak localization obtained from 50 water hammer experiments. The methods used for leak localization is (a) MFP-VE with all frequencies; (b) MFP-VE with only resonant frequencies; (c) MFP-E with all frequencies; (d) MFP-E with only resonant frequencies; (e) wavelet transform method; (f) FRF peak pattern method; (f) FRF peak pattern method; (g) ITA method. The ratio  $Q_0^L/Q_0$  of leak flow and main flow is 40%.

<i>e</i>	(a)	(b)	(c)	(d)	(e)	(f)	(g)
Mean [m]	0.52	17.15	5.05	74.16	17.60	14.10	12.15
Std. [m]	0.28	16.76	0.18	18.15	26.80	6.73	21.50
Max [m]	1.16	35.74	5.56	92.24	81.40	26.42	75.38
Min [m]	0.04	0.04	4.66	4.06	0.08	2.38	0.5
5% Percentile [m]	0.06	0.06	4.76	24.54	0.21	2.38	0.5
25% Percentile [m]	0.26	1.06	4.96	72.64	0.77	12.02	2.38
50% Percentile [m]	0.51	4.44	5.06	72.99	1.64	16.78	7.20
75% Percentile [m]	0.66	34.44	5.16	89.84	26.92	16.78	8.14
95% Percentile [m]	0.96	35.24	5.26	91.14	71.94	26.42	75.38

Table 5: Statistics of absolute error of leak localization obtained from 50 water hammer experiments. The methods used for leak localization is (a) MFP-VE with all frequencies; (b) MFP-VE with only resonant frequencies; (c) MFP-E with all frequencies; (d) MFP-E with only resonant frequencies; (e) wavelet transform method; (f) FRF peak pattern method; (g) ITA method. The ratio  $Q_0^L/Q_0$  of leak flow and main flow is 20%.

e	(a)	(b)	(c)	(d)	(e)	(f)	(g)
Mean [m]	0.56	21.80	4.94	79.30	18.85	14.29	23.94
Std. [m]	0.24	16.51	0.21	14.08	29.14	6.74	10.61
Max [m]	1.16	35.64	5.46	91.74	85.98	26.42	56.34
Min [m]	0.04	0.16	4.56	4.26	0.03	2.38	6.18
5% Percentile [m]	0.06	0.16	4.66	72.24	0.03	2.38	9.47
25% Percentile [m]	0.46	1.06	4.76	72.64	0.71	12.02	17.43
50% Percentile [m]	0.56	34.14	4.96	73.09	1.78	16.78	21.89
75% Percentile [m]	0.56	34.74	5.06	90.64	32.31	16.78	28.20
95% Percentile [m]	0.96	35.54	5.26	91.44	81.20	26.42	42.94

Table 6: Statistics of absolute error of leak localization obtained from 50 water hammer experiments. The methods used for leak localization is (a) MFP-VE with all frequencies; (b) MFP-VE with only resonant frequencies; (c) MFP-E with all frequencies; (d) MFP-E with only resonant frequencies; (e) wavelet transform method; (f) FRF peak pattern method; (g) ITA method. The ratio  $Q_0^L/Q_0$  of leak flow and main flow is 10%.

$Q_0^L/Q_0$	error	(a)	(b)	(c)	(d)	(e)	(f)	(g)
	$\leq 1 \text{ m}$	96%	32%	0	0	32%	0	6%
40%	$\leq 3 \text{ m}$	100%	44%	0	0	62%	28%	16%
	$\leq 5 \text{ m}$	100%	44%	24%	8%	76%	28%	30%
	$\leq 10 \text{ m}$	100%	44%	72%	10%	80%	28%	90%
	$\leq 1 \text{ m}$	96%	24%	0	0	38%	0	20%
20%	$\leq 3 \text{ m}$	100%	48%	0	0	60%	28%	28%
	$\leq 5 \text{ m}$	100%	50%	38%	4%	64%	20%	38%
	$\leq 10 \text{ m}$	100%	52%	100%	4%	66%	20%	90%
	$\leq 1 \text{ m}$	98%	22%	0	0	34%	0	0
10%	$\leq 3 \text{ m}$	100%	38%	0	0	58%	20%	0
	$\leq 5 \text{ m}$	100%	38%	64%	2%	64%	20%	0
	$\leq 10 \text{ m}$	100%	38%	100%	2%	68%	20%	16%

Table 7: Successful rate of leak localization computed from 50 water hammer experiments in the sense of absolute error less than 1 m, 3 m, 5 m and 10 m, respectively. The methods used for leak localization is (a) MFP-VE with all frequencies; (b) MFP-VE with only resonant frequencies; (c) MFP-E with all frequencies; (d) MFP-E with only resonant frequencies; (e) wavelet transaform method; (f) FRF peak pattern method; (g) ITA method.

# 427 5. Discussions

In this section, the experimental results and the leak localization methods are more profoundly discussed, in the following aspects.

- Modeling physical complexity: An advantage of the MFP approach pro-430 posed in this paper is that the algorithm includes the viscoelasticity ef-431 fect of pipe wall, which largely modifies the behavior of transient wave 432 in pipe and measured data. The physical model dependent methods 433 like MFP can produce very accurate leak localization. However, its 434 accuracy depends on the correctness of the physical model. When the 435 elastic (wrong) model is used in the viscoelastic case, i.e., the meth-436 ods (c) and (d), or the viscoelastic coefficients cannot be satisfactorily 437 given, the localization is not accurate. In the case that accurate model 438 is not available, methods less dependent on physical model, such as the 439 wavelet method (e), is an option. 440
- Sufficienct use of information from measurements: It has been shown 441 in the literature that many leak detection methods using partial in-442 formation (reflection, resonant frequencies, etc.) work well in ideal 443 environments. However, in real experiments which are usually very 444 noisy and have many unknown uncertainties, the used partial infor-445 mation is very possibly contaminated such that these methods fail. A 446 solution is to use as much information as possible. This paper exper-447 imentally justifies this issue: the method (a), which uses all available 448 information, outperforms the method (b), which uses only resonant 449 frequencies, and the method (e), which uses only the reflection due to 450 leak in the first period. This issue has also been justified numerically in 451 [21] and theoretically via the theory of Fisher information and Cramér-452 Rao lower bound (CRLB) [22, 51]: each time step of the time signal 453 or each frequency in the FRF are useful information for leak detection 454 and decreases the expected estimation error. 455
- Multi-sensor fusion: Again, due to the presence of noise and uncertainty, it is preferable to use as much information as possible. Except using more frequencies and time signals, another approach is to set more sensors in the pipe. MFP (as well as ITA) is able to fuse information from different sensors, without changing the algorithm but only increases the length of the vectors  $\Delta \mathbf{h}$  and  $\mathbf{G}$  in Eq. (36). On the

other hand, MFP needs at least three sensors, while the other three
methods (e), (f) and (g) can be used if only one sensor is available,
which provides more flexibility.

*Computation complexity*: MFP and ITA has the similarity that both 465 methods decide the leak by matching the measured data with the physi-466 cal model. However, ITA needs to solve a multiple-parameter optimiza-467 tion problem; its dimension equals to the number of potential leaks as-468 sumed in the pipe. Therefore, a high localization resolution or precision 469 of ITA implies an extremely slow computation and vice versa. Also, 470 the result of ITA strongly depends on the initial value and it is very 471 possible that the algorithm stops at local maximum which corresponds 472 to a bad result. By contrast, the matching of MFP only needs to solve 473 a 1D optimization, therefore a 1D search/plot of MFP objective func-474 tion versus leak location along the pipe is sufficient. This implies that 475 MFP has a fast computation, does not need initial guess of leak, and 476 avoids local maximum traps. 477

The above discussion regarding the properties of the leak localization methods is briefly summarized in Table 8.

	MFP	Wavelet	Peak pattern	ITA
Modeling viscoelasiticy	Yes	No	No	Yes
Sufficient use of signal	Yes	No	No	Yes
Initial parameter	No	No	No	Yes
Computation speed	Fast	Fast	Fast	Slow
Using multiple sensors	Yes	No	No	Yes
Minimum sensor number	3	1	1	1

Table 8: Summary of the properties of the leak localization methods.

# 480 6. Conclusion

This paper addresses the problem of leakage detection in a viscoelastic pipe. The viscoelasticity of pipe wall is modeled in the governing equations of transient wave and, in the frequency domain, the viscoelastic effect is equivalent to modifying the transfer matrix with a frequency-dependent wave speed. Then, Matched-field processing (MFP), which has been proposed for leakage detection in elastic pipes, is generalized for the viscoelastic case.

Transient experiment results with viscoelastic pipe in the Water Engi-487 neering Laboratory at University of Perugia and in the Water Resources 488 Research Laboratory at Hong Kong University of Science and Technology 489 are both studied to justify the proposed method. It is shown that both con-490 sidering pipe wall viscoelasticity and using more frequencies (instead of using 491 only the resonant frequencies) in the leakage localization scheme are essential 492 for leak localization accuracy. The proposed MFP method outperforms clas-493 sical leakage detection methods in the literature in the sense of smaller leak 494 localization error. For a small leak where the flow ratio of leak and main pipe 495 is 10%, MFP is very accurate that, in 49 of 50 experiments, the leak local-496 ization error is lower than 1 m (the error of the rest one is 1.14 m), although 497 some of these signals are seriously contaminated by noise and uncertainties. 498 Further work may be conducted in several directions. First, the gener-499 alization of the results to more complicated cases, for example localization 500 of multiple leaks and application to more complex piping systems, would 501 be important. Besides, this paper uses the pipe wall viscoelastic coefficients 502 obtained from a no-leak test, which may not be available in some practical 503 cases. Therefore, a more advanced method that can jointly estimate leaks 504 and the viscoelastic coefficients would be interesting. 505

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